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# High mobility field-effect transistors based on MoS<sub>2</sub> crystals grown by the flux method

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## Abstract

Two-dimensional (2D) molybdenum disulphide (MoS<sub>2</sub>) transition metal dichalcogenides (TMDs) have great potential for use in optical and electronic device applications; however, the performance of MoS<sub>2</sub> is limited by its crystal quality, which serves as a measure of the defects and grain boundaries in the grown material. Therefore, the high-quality growth of MoS<sub>2</sub> crystals continues to be a critical issue. In this context, we propose the formation of high-quality MoS<sub>2</sub> crystals via the flux method. The resulting electrical properties demonstrate the significant impact of crystal morphology on the performance of MoS<sub>2</sub> field-effect transistors. MoS<sub>2</sub> made with a relatively higher concentration of sulphur (a molar ratio of 2.2) and at a cooling rate of 2.5 °C h<sup>-1</sup> yielded good quality and optimally sized crystals. The room-temperature and low-temperature (77 K) electrical transport properties of MoS<sub>2</sub> field-effect transistors (FETs) were studied in detail, with and without the use of a hexagonal boron nitride (h-BN) dielectric to address the mobility degradation issue due to scattering at the SiO<sub>2</sub>/2D material interface. A maximum field-effect mobility of 113 cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup> was achieved at 77 K for the MoS<sub>2</sub>/h-BN FET following high-quality crystal formation by the flux method. Our results confirm the achievement of large-scale high-quality crystal growth with reduced defect density using the flux method and are key to achieving higher mobility in MoS<sub>2</sub> FET devices in parallel with commercially accessible MoS<sub>2</sub> crystals.

Keywords: MoS<sub>2</sub>, Flux method, Field-effect transistor, hexagonal boron nitride (h-BN)

(Some figures may appear in colour only in the online journal)

## 1. Introduction

Over the last two decades, transition metal dichalcogenides (TMDs) have been a focal point, owing to their potential application in next-generation electronic and optoelectronic

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**Table 1.** Process parameters of high-quality crystal MoS<sub>2</sub> growth via the flux method.

Material	Purity	Material Molar Ratio			
		Cooling rate: 2.5 °C h <sup>-1</sup>		Cooling rate: 2 °C h <sup>-1</sup>	
		Sample #1	Sample #2	Sample #3	Sample #4
Mo	4N	1	1	1	1
S	5N	2	2.2	2	2.2
Sn	5N	15	15	15	15

devices. One of these TMDs is 2D MoS<sub>2</sub>, which consists of covalently coupled S-Mo-S (sulphur-molybdenum-sulphur) sheets with weak van der Waals forces between them. The MoS<sub>2</sub> bandgap ranges from 1.2 eV (bulk) to 1.8 eV (single layer), depending on the number of coupled layers, allowing different applications [1, 2]. However, certain challenges need to be addressed before MoS<sub>2</sub> can be used in actual device applications. One of these challenges is the crystal quality of grown MoS<sub>2</sub>, which is an important factor in the production of electronic devices with respect to performance [3, 4]. The crystal quality, number of layers, and the chemical composition are factors that could degrade the fabricated device performance. Therefore, crystal growth techniques play an important role in determining the crystal quality and hence the device performance [5].

Good quality MoS<sub>2</sub> crystals can be obtained via vapor deposition techniques such as chemical vapor deposition (CVD), chemical vapor transport (CVT), pulsed laser deposition (PLD), or other chemical synthesis routes. There are many reports of the successful growth of MoS<sub>2</sub> crystals by CVD, CVT, and PLD [6–9]. The crystal quality appears to be dependent on parameters such as growth temperature, precursor concentration, gas-flow rates, and the growth pressure of the furnace. Regardless of the success of CVD, issues such as tensile stress due to substrate interaction and by-products of a possible loss of reaction result in impurities in the grown MoS<sub>2</sub> layers. This leads to disadvantages such as cracks and grain boundaries in the material, resulting in intrinsic rippling in the layered regions [10–13]. Recently, the flux method has gained attention for yielding high-quality, defect-free, large-area 2D crystals. The flux method is a solution-based growth technique, whereby the required material is dissolved in a solvent (flux). The grown crystals are free from thermal strain, and the desired quality of crystals can be obtained by optimizing the process temperature. 2D crystals grown by the flux method possess a defect density that is one to two orders lower than that of commercially available crystal materials [14]. Zhang *et al* [15] reported the growth of MoS<sub>2</sub> crystals using a tin (Sn) flux and explored its use for photonic applications. Cevallos *et al* [16] used a common salt flux to grow MoS<sub>2</sub> and WS<sub>2</sub> layered crystals and obtained excellent photoluminescent properties. Furthermore, flux-grown MoS<sub>2</sub> crystals can also be utilised for solar cells [17], photoluminescence [18], and various frequency-generation [19] applications.

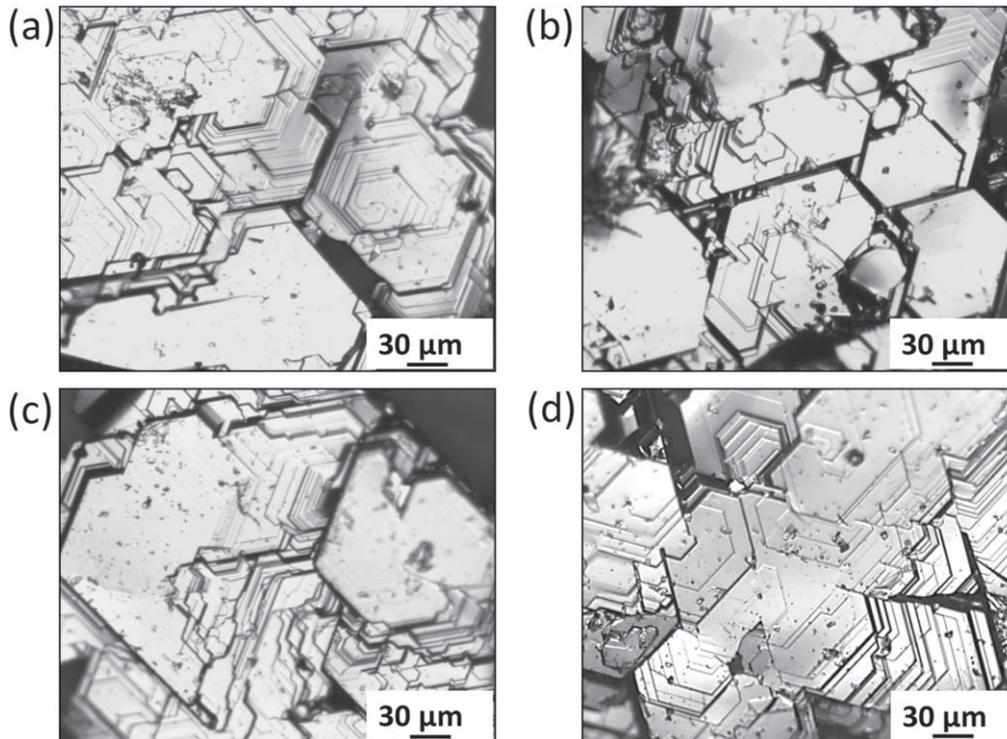
Although the MoS<sub>2</sub> crystals grown via the flux method retain excellent optical properties, the desired electrical

properties have not yet been studied. In addition, the 2D material properties are highly reliant on the supporting dielectric substrate used in the device fabrication. SiO<sub>2</sub> has been a popular choice of dielectric for device fabrication, due to its ideal interface properties when used with Si. However, 2D layers exfoliated on SiO<sub>2</sub> tend to suffer from mobility degradation due to carrier scattering and trapped charges at the SiO<sub>2</sub>/2D layer interface [20]. To address this problem, the use of hexagonal boron nitride (h-BN) layers has been proposed by various groups in recent studies [21, 22]. An h-BN substrate is an excellent support for the exfoliated MoS<sub>2</sub>, improving device performance in terms of mobility. In this study, we report the electrical properties of high-quality MoS<sub>2</sub> crystals grown using the flux method and mechanically exfoliated for device applications with and without h-BN. This study aims to elucidate the relationship between the cooling rate and molar ratio in MoS<sub>2</sub> crystals and their electrical mobility. This study reveals that the cooling rate has a significant impact on the material's mobility, which is strongly related to crystal size and quality. We prepared four different ampules: the first two ampules had a 1:2:15 molar ratio and the other two had 1:2.2:15 molar ratios. Afterwards, two different molar ratio ampules were prepared at a 2.5 °C h<sup>-1</sup> cooling rate and the other two were prepared at 2.0 °C h<sup>-1</sup>. The grown MoS<sub>2</sub> flux was characterized by x-ray diffraction (XRD) and Raman spectroscopy. The high-quality flux-grown MoS<sub>2</sub> crystals mechanically exfoliated on an h-BN dielectric showed great improvement in mobility, compared to that of a non-h-BN layer device measured at 10<sup>-3</sup> Torr of pressure.

## 2. Experimental details

### 2.1. MoS<sub>2</sub> crystal growth by flux method

To grow high-quality crystals, the method reported by Zhang *et al* [15] was adopted. Here, we used a Sn flux, and Sn did not contribute to the reaction because it is an unreactive solvent. It has a lower melting point than sulphur (S) and molybdenum (Mo), and it also remains stable during the high-temperature process allowing the formation of MoS<sub>2</sub>, preventing impurities from being added. It is also easy to remove or clean after the formation of MoS<sub>2</sub> crystals [16]. The four different samples were prepared with two different molar ratios and different cooling rates. Table 1 gives the details for all four samples. To check whether it was possible to



**Figure 1.** Microscopic images of MoS<sub>2</sub> crystals grown by the flux method (a) with a molar ratio of 1:2:15 and a cooling rate of 2.5 °C h<sup>-1</sup> (b) a molar ratio of 1:2.2:15 and a cooling rate of 2.5 °C h<sup>-1</sup> (c) a molar ratio of 1:2:15 and a cooling rate of 2 °C h<sup>-1</sup> (d) a molar ratio of 1:2.2:15 and a cooling rate of 2 °C h<sup>-1</sup>. Image (b) shows the clear formation of large crystals.

overcome sulphur vacancies in the MoS<sub>2</sub> crystals, one sample was prepared with a high sulphur concentration.

The source materials were added into an alumina crucible fixed into a quartz tube ampule, which was later vacuum sealed and loaded into a high-temperature furnace at 1100 °C at a ramp-up rate of 50 °C h<sup>-1</sup>. The temperature was held constant for 48 h and then the samples were allowed to cool down at 2.5 °C h<sup>-1</sup> and 2 °C h<sup>-1</sup>, respectively. The resultant ampules were centrifuged at 2000 rpm for 60 s. The grown crystals were characterized by Raman spectroscopy and x-ray diffraction techniques to verify the crystal formation and crystal quality. The formed crystals were then used for device fabrication.

## 2.2. Device fabrication and characterization

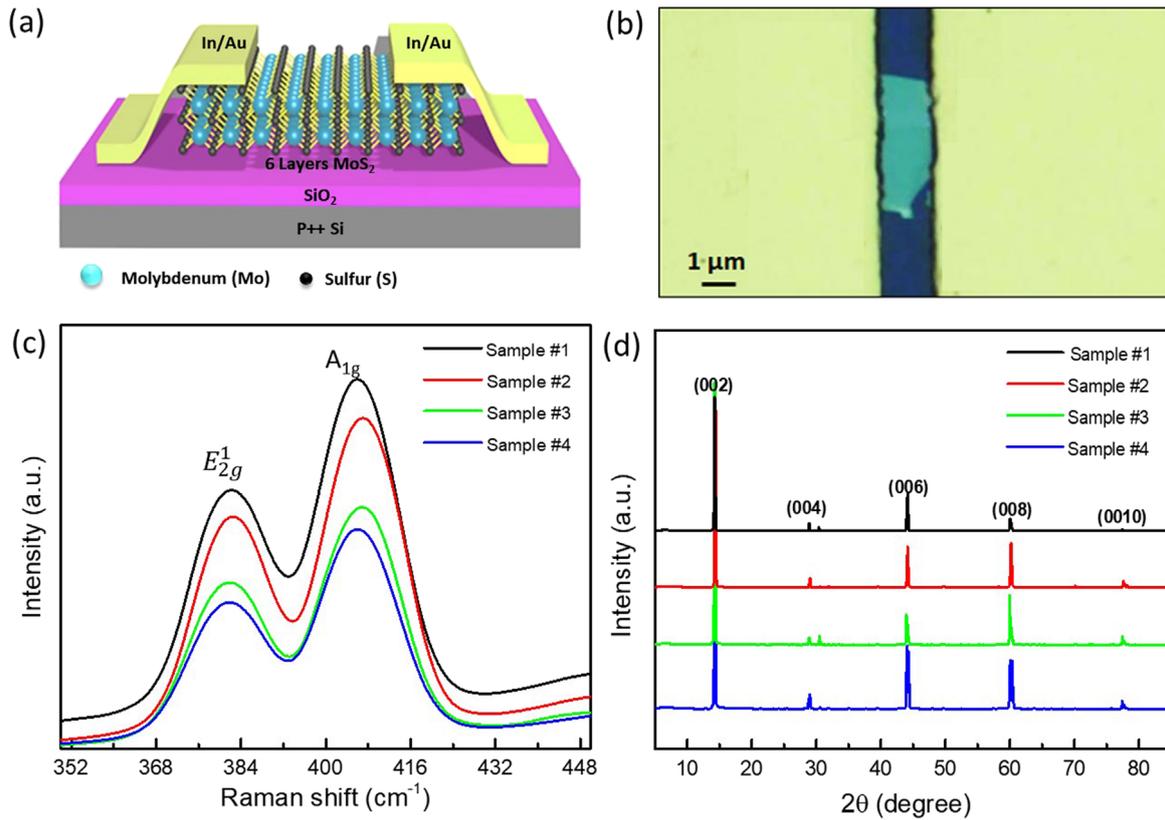
Four different devices with the same number of 2D layers (six layers) were fabricated via mechanical exfoliation of the flux-grown crystal to obtain MoS<sub>2</sub> crystals on a (285 nm) SiO<sub>2</sub>/Si substrate. The substrate was pre-cleaned with acetone and isopropyl alcohol (IPA) before mechanical exfoliation. Exfoliated MoS<sub>2</sub> flakes were then transferred onto the h-BN/SiO<sub>2</sub>/Si substrate using the dry transfer polydimethylsiloxane, and polycarbonate (PDMS and PC stamp) method. The source and drain electrodes were patterned through a standard photolithographic process. Au/In (10/30 nm) bilayer metal electrodes were deposited using the e-beam evaporation technique. A lift-off process was carried out to remove unwanted metal and photoresist from the sample. Finally, the device was subjected to post-metallization

annealing at 120 °C for 90 min in an argon gas atmosphere to enhance the contact properties. Figure 2(a) shows a typical schematic of the fabricated device structure. After verifying the crystal quality of all four samples, a device with a few layers of h-BN used as a dielectric between the SiO<sub>2</sub> and the MoS<sub>2</sub> was fabricated, using similar process parameters for fabrication to those given above.

The flux-grown MoS<sub>2</sub> crystals exfoliated onto Si/SiO<sub>2</sub> substrates were visualised using an Olympus BX51 optical microscope at a resolution of 20 μm. A Renishaw confocal Raman spectrometer with a monochromatic laser source with a wavelength of 514.5 nm was used to characterize the molecular vibrations in the six-layer MoS<sub>2</sub> flake. The crystal quality was examined using a Bruker Advance D8 x-ray diffractometer at various angular settings between 5° and 85°. The electrical characteristics such as the I<sub>ds</sub>-V<sub>gs</sub> and I<sub>ds</sub>-V<sub>ds</sub> of all five devices with and without the h-BN were investigated by measuring the device at room temperature (RT) and at 77 K using a Keithley-4200 interactive test module. All the electrical measurements were performed in a vacuum environment at a pressure of 2 × 10<sup>-3</sup> Torr.

## 3. Results and discussion

Optical images of MoS<sub>2</sub> crystals grown under different conditions are shown in figures 1(a)–(d). Figure 1 intimates that the cooling rate has a significant impact on crystal morphology.



**Figure 2.** (a) A 3D schematic of the six-layer MoS<sub>2</sub> device with In/Au metal contacts fabricated on 285 nm of SiO<sub>2</sub>/Si. (b) Optical image of the real device. (c) Raman spectrum of six-layer MoS<sub>2</sub> crystals grown under four different sets of conditions via the flux method at room temperature. (d) XRD pattern of a 2H-MoS<sub>2</sub> high-quality single crystal formed under four different sets of conditions. Sample #1. 1:2:15, 2.5 °C h<sup>-1</sup>; sample #2. 1:2.2:15, 2.5 °C h<sup>-1</sup>; sample #3. 1:2:15, 2 °C h<sup>-1</sup>; and sample #4. 1:2.2:15, 2 °C h<sup>-1</sup>.

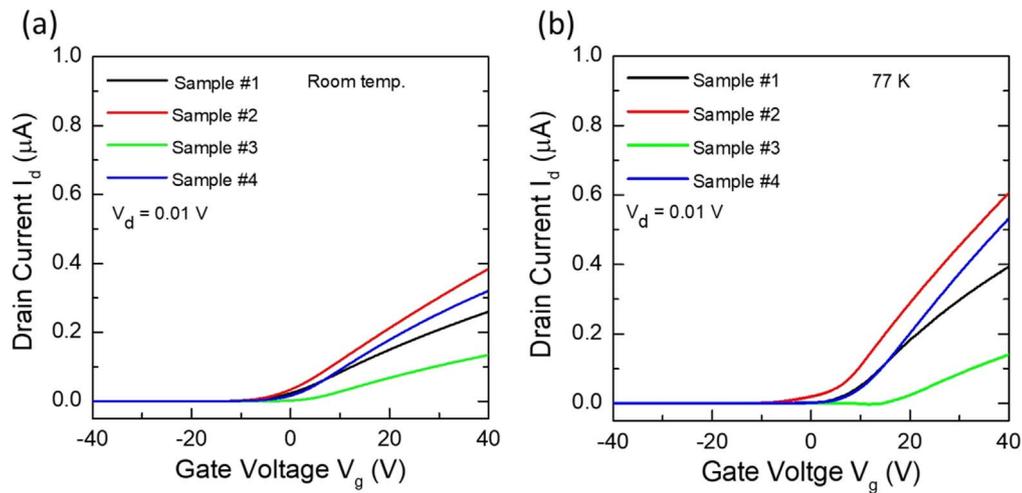
Figure 1(a) is an optical image of the crystal formed at a 2.5 °C h<sup>-1</sup> cooling rate with a sulphur to molybdenum ratio of 2:1. Cracks in the crystal were observed at lower molar sulphur concentrations. Figure 1(b) shows high-quality MoS<sub>2</sub> single crystals in the sample for which the cooling rate was maintained at 2.5 °C h<sup>-1</sup> and which had 1:2.2:15 molar ratios of Mo, S, and Sn, respectively. Furthermore, the cooling rate was decreased to 2.0 °C h<sup>-1</sup> and with similar sulphur concentrations, as above. Figure 1(c) shows that lower-quality crystal formation tends towards the polycrystalline phase with 1:2:15 molar ratios. We further increased the sulphur molar ratio to 2.2, resulting in good quality crystals at a cooling rate of 2.0 °C h<sup>-1</sup>, as shown in figure 1(d). The optical images show the formation of crystalline MoS<sub>2</sub> crystals grown in the Sn flux with a greater decrease in the cooling rate. The crystals obtained were later used to investigate the effects of crystal morphology on the optical and electrical properties of MoS<sub>2</sub> FETs.

Raman spectroscopy was used to verify the formation of six-layer single-crystal MoS<sub>2</sub> over SiO<sub>2</sub>/Si. A high-resolution Raman microscope was used to measure the Raman shift in the grown samples at an excitation wavelength ( $\lambda_{ex}$ ) of 532 nm. Figure 2(c) shows the Raman spectra for all four samples grown by the flux method. The two sharp peaks at 406.8 cm<sup>-1</sup> and 382.18 cm<sup>-1</sup> correspond to the A<sub>1g</sub> active mode and the E<sub>2g</sub><sup>1</sup> vibrational modes. The A<sub>1g</sub> mode is related

to the in-plane vibration of the S atoms, while E<sub>2g</sub><sup>1</sup> suggests the out-of-plane vibration of the Mo and S atoms. The peak separation of 24.5 cm<sup>-1</sup> is the same for all samples, demonstrating that the number of layers is six for each device [23]. Furthermore, the peak intensity is observed to decrease as the crystal size increases, indicating the large crystal formation of sample 4. The Raman peak intensity is also used to demonstrate the surface adsorptions and surface functionalization conditions describing the optical transitions at different excitation energies [24, 25]. The peak intensity is dependent on the excitation energy used during measurement. A higher excitation energy can lead to a higher peak intensity; however, the E<sub>2g</sub> mode always has a lower intensity than the A<sub>1g</sub> mode, as it is dependent on the A<sub>1g</sub> close to the C excitation energy band, which is found to be similar here for the case of MoS<sub>2</sub> crystals grown by the flux method.

Furthermore, the decrease in the A<sub>1g</sub> and E<sub>2g</sub><sup>1</sup> peak intensities attributed to the structural changes between the interlayers, the lattice vibration due to variation in the cooling rate and the sulphur content during the growth of MoS<sub>2</sub> flakes, and the mechanical exfoliation would not cause the change in the intensity of the Raman peak.

The x-ray diffraction (XRD) pattern in figure 2(d) shows single-crystal 2H MoS<sub>2</sub> peaks for all four samples. The 2θ peaks at 14.2°, 29.1°, 44.2°, 60.2°, and 77.5° correspond to the (002), (004), (006), (008), and (0010) planes of the MoS<sub>2</sub>



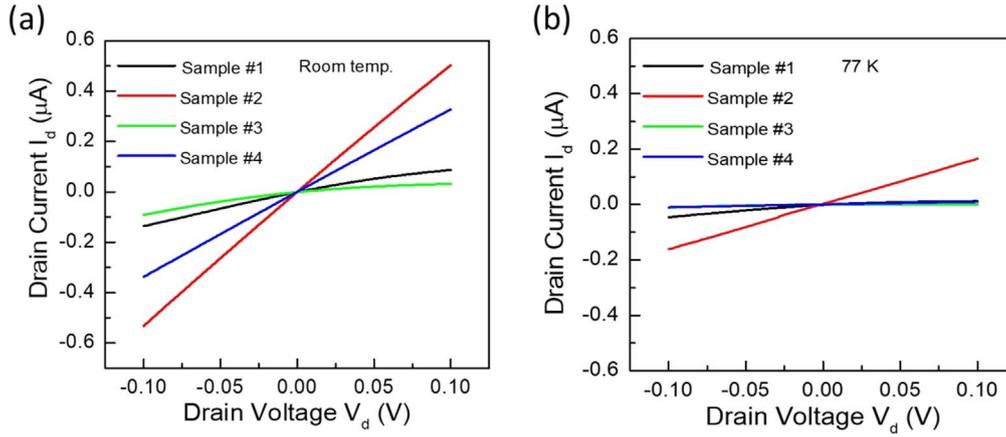
**Figure 3.** Electrical properties of the MoS<sub>2</sub> FET. (a) Transfer characteristics of the four different devices measured at room temperature with  $V_{ds} = 0.01$  V and back-gate voltages varying from  $-40$  V to  $40$  V. (b) Low-temperature transfer characteristics of the MoS<sub>2</sub> FET measured at  $77$  K with  $V_{ds} = 0.01$  V and gate voltages ranging  $-40$  to  $40$  V. All measurements were carried out at  $2 \times 10^{-3}$  Torr pressure.

hexagonal structure, respectively, and these results are consistent with those of the Joint Committee on Powder Diffraction Standards (JCPDS) card no. 37-1492. The (001) diffraction plane, attributed to a d-spacing of  $0.62$  nm, was observed for a single-crystal MoS<sub>2</sub> flake [26]. Further, the additional peak at approximately  $30.59^\circ$  in samples 1 and 3 is related to the presence of defects, while the shift to lower angles for the sulphur peak at  $31.40^\circ$  gives an increase in interlayer spacing from  $0.28$  to  $0.29$  nm [27]. However, this peak is insignificant in the XRD spectra of samples 2 and 4 because of the additional sulphur molar concentration, which therefore helps to overcome these defects. In the crystal structure of materials, the ab-plane strongly shows the two-dimensional growth behaviour of MoS<sub>2</sub> for samples 2 and 4, matching the similar behaviour for high-quality crystal growth results reported by Zhang *et al* [15]. Therefore, we can consider the MoS<sub>2</sub> crystals grown via the Sn flux method described in this study to be high-quality single crystals.

Figure 3 shows the transfer characteristics ( $I_{ds}$  versus  $V_{gs}$ ) of the non-h-BN MoS<sub>2</sub> field-effect transistors measured at RT (a) and  $77$  K (b) at a small applied drain voltage ( $V_{ds}$ ) of  $0.01$  V for four different crystals. All four devices show a linear increment in the drain current with respect to the gate voltage, and n-type conductive behaviour. Among them, the devices with higher sulphur concentrations of  $2.2$  M exhibit an improvement in the drain current, as observed in figures 3(a) and (b) measured at RT and  $77$  K, respectively, except for sample 3 at  $77$  K, which shows a reduction in the drain current. The possibility is undeniable that sulphur vacancies are present, which later become a source of grain boundaries in the exfoliated samples, as reported by various groups through scanning tunnelling microscopy [14]. The introduction of O<sub>2</sub> or H<sub>2</sub> annealing can heal this vacancy, as reported earlier [28–30].

However, the extra sulphur concentration added during the MoS<sub>2</sub> crystal formation may occupy these vacancies during exfoliation and improve the electrical performance of the MoS<sub>2</sub> FET, as can be observed from the transfer

characteristics of the RT and  $77$  K samples. In relation to the crystal quality, samples 2 and 4 show high crystal quality, despite the large crystal size of the samples. The low cooling rates of  $2^\circ\text{C h}^{-1}$  could make the material of sample 3 polycrystalline [15], which degrades the drain current of the device, but a high concentration of sulphur helps to heal these vacancies; thus, the current in sample 4 is higher than that in sample 3 but lower than that in sample 2. All four samples show an  $I_{on/off}$  ratio of  $10^6$ . A negative threshold voltage shift of  $-4.6$  V is observed in sample 2. The transfer characteristics measured at  $77$  K (figure 3(b)) show thermally activated charge transport in all four devices, with a similar trend observed for the RT measured devices. This thermally activated charge transport is attributed to the hopping mechanism caused by defect-induced localized states, where the carriers are pinned within the specialised locations of these disorders which are either vacancies or some interstitial states. However, pinned carriers located near the Fermi level gain thermal energy and hop to their nearest neighbours by thermally activated hopping, or hop further to other locations, subject to an energy minimization constraint, in the variable-range hopping mode. This causes a negative temperature dependence of the mobility. The observed phenomenon is similar to the results reported by Ghatak *et al* for MoS<sub>2</sub> grown on SiO<sub>2</sub> [31]. Furthermore, this also can be realised in terms of a Schottky transition within the device due to the grain boundaries; however, the Schottky transition is quite untraceable in these devices, except for sample two, as depicted in figure 3(b). It is also noted that thermally activated trap charges also cause a shift in the threshold voltage of multilayer MoS<sub>2</sub> devices [32]. The threshold voltage shift and change in contact resistance are both visible in figure 3(b), which confirms the presence of thermally activated transport in all four devices. The threshold voltage increases towards the positive side with an increase in the drain current for all four samples. The field-effect mobilities for devices measured at RT and  $77$  K were extracted from the measured



**Figure 4.** Output characteristics of the MoS<sub>2</sub> FET. (a)  $I_{ds}$  versus  $V_{ds}$  curve measured at room temperature for four different samples. (b) The  $I_{ds}$  versus  $V_{ds}$  curve measured at 77 K. Both measurements were carried out with a zero gate voltage. All measurements were carried out at  $2 \times 10^{-3}$  Torr of pressure.

**Table 2.** Comparisons of the electrical parameters for samples 1 to 4. The six-day growth process uses the  $2.5 \text{ }^\circ\text{C h}^{-1}$  cooling rate, while the eight-day process uses the  $2.0 \text{ }^\circ\text{C h}^{-1}$  cooling rate.

Sample name	Molecular ratio (Mo:S:Sn)	$I_{on/off}$ ratio	Threshold voltage ( $V_{th}$ ) in (V)	Mobility ( $\mu$ ) in ( $\text{cm}^2 \text{V}^{-1} \text{s}^{-1}$ )
Room temperature				
Sample #1	1:2:15	$10^6$	-4.3	27.3
Sample #2	1:2.2:15	$10^6$	-4.6	45.1
Sample #3	1:2:15	$10^6$	2.3	16.5
Sample #4	1:2.2:15	$10^6$	-0.9	44.5
77 K				
Sample #1	1:2:15	$10^6$	5.4	54.2
Sample #2	1:2.2:15	$10^6$	3.4	88.4
Sample #3	1:2:15	$10^6$	13.4	25.0
Sample #4	1:2.2:15	$10^6$	7.9	86.9

conductance ( $g_{max}$ ) using the following equation [30]:

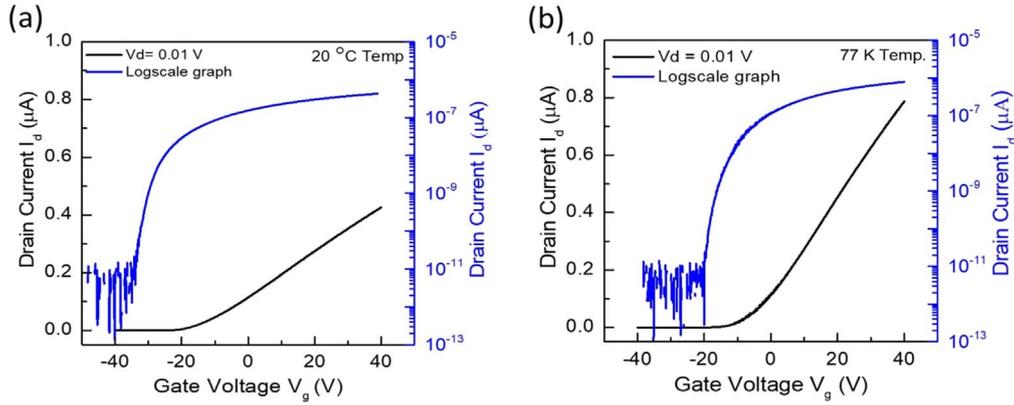
$$\mu = \left( \frac{L}{W} \frac{g_{max}}{C_{ox} V_{ds}} \right) \quad (1)$$

where  $g_{max}$  is the maximum peak value of  $dI_{ds}/dV_{gs}$ . A maximum room temperature mobility of  $45.1 \text{ cm}^2 \text{V}^{-1} \text{s}^{-1}$  was calculated for sample 2; while a field-effect mobility of  $88.4 \text{ cm}^2 \text{V}^{-1} \text{s}^{-1}$  was calculated for sample 2 at 77 K. The observed mobility is in good agreement with the mobility reported by Lin *et al* [33] for a five-layer MoS<sub>2</sub> ( $110 \text{ cm}^2 \text{V}^{-1} \text{s}^{-1}$ ) FET non-h-BN device at 77 K measured at  $10^{-6}$  Torr pressure. It is worth noting that the field-effect mobility is sensitive to the external environment. In this study, all measurements at 77 K were carried out at a  $2 \times 10^{-3}$  Torr of base pressure. Therefore, the extracted mobility is highest at this pressure at RT and 77 K, which to the best of our knowledge, reflects the best quality single-crystal formation of sample 2 with a  $2.5 \text{ }^\circ\text{C h}^{-1}$  cooling rate. Unlike a monolayer MoS<sub>2</sub>, the increase in field-effect mobility at low temperature is also assisted by the reduction in phonon scattering at low temperatures, because the multilayered MoS<sub>2</sub> favours weaker phonon-electron interactions due to reduced lattice vibrations [34]. Table 2 shows the

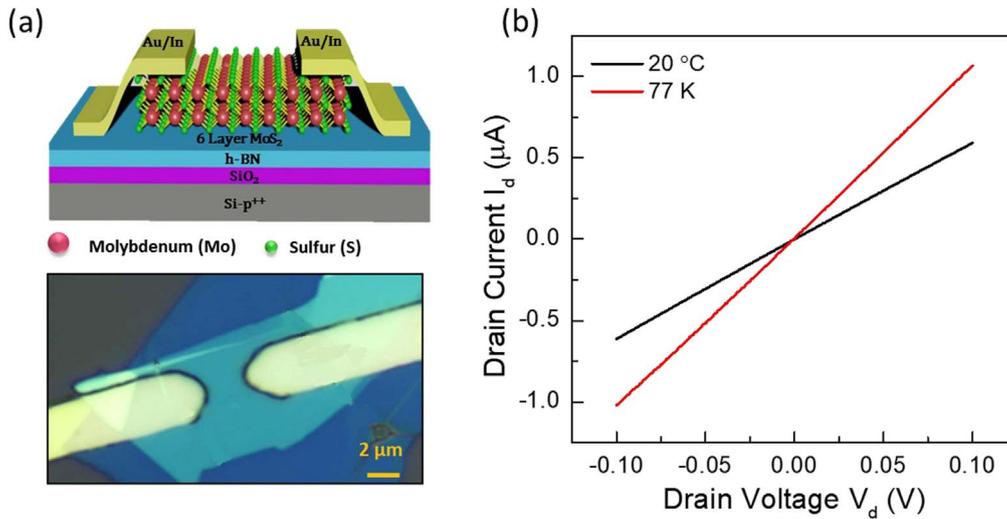
summary of all the calculated electrical parameters at RT and 77 K for MoS<sub>2</sub> FET devices grown via the flux method. Figures 4(a) and (b) show the output characteristics ( $I_{ds}$  versus  $V_{ds}$ ) with no applied back-gate voltage ( $V_{gs}$ ) for samples 1–4, characterized at room temperature and 77 K, respectively.

The linear behaviour with increasing applied drain voltage reflects the good ohmic contact between the exfoliated MoS<sub>2</sub> and the Au/In metal contact device and shows the strong dependence of the output drain current on the crystal morphology of the fabricated devices. The highest  $I_{ds}$  values were observed in samples 2 and 4 in figure 4, which displayed a strong dependency on the crystal quality similar to that discussed previously. Meanwhile, the output characteristics measured at 77 K in figure 4(b) show a deviation from the linear behaviour and a significant reduction in the drain current for all three samples, with the exception of sample 2.

Further, a transition from ohmic behavior for a Schottky barrier was observed for lower temperatures, which can be attributed to the grain boundary impurities in the samples grown under different conditions. These grain boundaries allow for a metal charge injection within the channel material at lower temperatures [35]. However, this metal injection can be improved at lower pressures close to approximately



**Figure 5.** Electrical properties of the MoS<sub>2</sub>/h-BN FET. (a)  $I_d$ - $V_g$  characteristics at room temperature with  $V_{ds} = 0.01$  V and back-gate voltages varying from  $-40$  V to  $40$  V. (b)  $I_d$ - $V_g$  characteristics at  $77$  K with  $V_{ds} = 0.01$  V and gate voltages ranging from  $-40$  V to  $40$  V. All measurements were carried out at  $2 \times 10^{-3}$  Torr of pressure.



**Figure 6.** (a) Shows the device schematic and an optical device image of the real device. (b)  $I_{ds}$  versus  $V_{ds}$  curves of the MoS<sub>2</sub>/h-BN FET measured at room temperature and  $77$  K. Both measurements were carried out at a zero gate voltage. All measurements were carried out at  $2 \times 10^{-3}$  Torr of pressure.

**Table 3.** A comparison of the electrical parameters for samples 1 to 4.

Device structure	Measurement conditions	$I_{on}/I_{off}$ ratio	Threshold voltage (V)	Mobility $\text{cm}^2 \text{V}^{-1} \text{s}^{-1}$
MoS <sub>2</sub>	RT	$10^6$	$-4.6$	$45.1$
	$77$ K	$10^6$	$3.4$	$88.4$
MoS <sub>2</sub> /h-BN	RT	$10^5$	$-23.16$	$50.5$
	$77$ K	$10^5$	$-10.54$	$113$

$10^{-6}$  Torr or by encapsulating the MoS<sub>2</sub> channel material using high- $k$  dielectrics such as HfO<sub>2</sub> [36].

The transfer characteristics when h-BN is inserted as a dielectric layer between the SiO<sub>2</sub> and MoS<sub>2</sub> layers show high mobility (sample 2) among the four samples, as shown in figure 5. Figure 5(a) shows the measurements performed at RT and the  $77$  K measurements are shown in figure 5(b). As expected, both the room temperature and  $77$  K measurements show an improvement in the drain current with the insertion of few layers of h-BN between the MoS<sub>2</sub> and the SiO<sub>2</sub>, compared to that of the non-h-BN MoS<sub>2</sub>/SiO<sub>2</sub> FET. The

insertion of the h-BN helps to improve the interface quality of 2D TMDs and SiO<sub>2</sub>, and this improvement therefore enhanced the electric field of the device and the drain current.

Figure 6(a) shows a device schematic and an optical device image of the MoS<sub>2</sub>/h-BN device; figure (b) shows the output characteristics of the MoS<sub>2</sub>/h-BN FET device measured at RT and  $77$  K.

The hexagonal boron nitride (h-BN) layer is less rough than the SiO<sub>2</sub> surface and has no surface doping effect, resulting in improved mobility. The field-effect mobility was extracted for devices using equation (1). The device measured

at RT had a mobility of  $50.5 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ , while the highest mobility of  $113 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  was achieved at 77 K at a pressure of  $2 \times 10^{-3}$  torr. Table 3 compares the electrical properties of the MoS<sub>2</sub> FETs with and without h-BN.

#### 4. Conclusions

The effect of crystal morphology on the electrical properties of six-layer MoS<sub>2</sub> FETs was investigated. The Sn flux method was adopted to grow high-quality single-crystal MoS<sub>2</sub> material. The optical images confirm the variations in the crystal dimensions for different cooling rates, while the Mo and S molar ratio helps to modify the crystallinity. The Raman shift was used to verify the number of layers of exfoliated MoS<sub>2</sub> on SiO<sub>2</sub>/Si. Room-temperature and low-temperature (77 K) measurements were carried out to investigate the electrical properties of the fabricated FET devices. The electrical properties revealed the great impact of crystal morphology on device performance. High-quality crystals with small grain boundaries and defects were obtained at a higher sulphur concentration and a cooling rate of  $2.5 \text{ }^\circ\text{C h}^{-1}$ , showing an improvement in the field-effect mobility. Furthermore, the addition of few layers of h-BN improves the mobility significantly, as discussed. Our study shows that the flux method is effective in achieving crystal growth; in addition, the crystal quality of MoS<sub>2</sub> can be tuned by varying the molar ratio between Mo and S and the cooling rate during the growth process. Moreover, high-quality crystals with the addition of an h-BN dielectric could be the key to achieving higher mobility in MoS<sub>2</sub>-based FET devices.

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#### Data availability statement

The data that support the findings of this study are available upon reasonable request from the authors.

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